Designability of lattice model heteropolymers

G. Tiana,¹ R. A. Broglia,^{1,2} and D. Provasi¹

¹Dipartimento di Fisica, Universitá di Milano e INFN Sezione di Milano, via Celoria 16, 20133 Milano, Italy

²The Niels Bohr Institute, Blegdamsvej 17, 2100 Copenhagen, Denmark

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Protein folds are highly designable, in the sense that many sequences fold to the same conformation. In the present work we derive an expression for the designability in a 20-letter lattice model of proteins which, relying only on the central limit theorem, has a generality which goes beyond the simple model used in its derivation. This expression displays an exponential dependence on the energy of the optimal sequence folding on the given conformation measured with respect to the lowest energy of the conformational dissimilar structures, an energy difference which constitutes the only parameter controlling designability. Accordingly, the designability of a native conformation is intimately connected to the stability of the sequences folding to them.

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I. INTRODUCTION

Even a quick look at the set of known proteins (protein data bank database) reveals a striking feature. While there are tens of thousands of protein sequences, they only assume some thousands of folds. In other words, proteins are highly designable. This concept can be quantified by measuring the number of sequences that fold uniquely into a particular structure.

With the use of a simple 20-letter lattice model [1-4] of protein folding it has been shown [5] that the whole issue of estimating the number *n* of sequences that fold to the same conformation is reduced to enumerate how many of them have native energy lying below a threshold E_c , the energy which any sequence with the same composition displays in the conformation structurally dissimilar to the native conformation [2].

The aim of the present paper is to provide a reliable, analytic expression for n, which we shall show increases exponentially with the gap δ between the native energy of the optimal sequence E_n and the threshold energy E_c . This functional form is found to be universal, as it emerges from the central limit theorem. We have, furthermore, found that while the parameters defining n depend on the interaction matrix, they are independent of the particular choice made for the native structure or for the optimal sequence.

In Secs. II and III we briefly review the 20 letters lattice model of proteins in general and the question of protein designability in particular. The quantitative, analytic answer to the question of how many mutations a designed protein tolerates is given in Sec. IV. The conclusions are collected in Sec. V.

II. LATTICE MODELS

A useful theoretical approach to study protein folding is provided by a simplified lattice model, where the protein is a string of beads that is arranged on a cubic lattice [6-8]. The configurational energy of a chain of N monomers is given by

$$E = \frac{1}{2} \sum_{i,j}^{N} U_{m(i),m(j)} \Delta(|\vec{r}_i - \vec{r}_j|), \qquad (1)$$

where $U_{m(i),m(j)}$ is the effective interaction potential between monomers m(i) and m(j), $\vec{r_i}$ and $\vec{r_j}$ denote their lattice positions, and $\Delta(x)$ is the contact function. In Eq. (1) the pairwise interaction is different from zero when *i* and *j* occupy nearest-neighbor sites, i.e., $\Delta(a)=1$ and $\Delta(na)=0$ for $n \ge 2$, where *a* indicates the step length of the lattice. In addition to these interactions, it is assumed that on-site repulsive forces prevent two amino acids from occupying the same site simultaneously, so that $\Delta(0)=\infty$.

We shall consider throughout a 20-letter representation of protein sequences where U is a 20×20 matrix. A possible realization of this matrix is given in Ref. [9] (Tables 5 and 6) where it was derived from frequencies of contacts between different amino acids in protein structures. The employment of a 20 \times 20 matrix ensures that the threshold energy E_c is well defined, depending only on the interaction matrix elements and on the composition of the protein in terms of amino acids. The model we study here is a generic heteropolymer model that has been shown to reproduce important generic features of protein-folding thermodynamics and kinetics, independent on the particular potential chosen [10,11]. However, in using such an approach, one should keep in mind that the labeling of amino acids (spherical beads all of the same size and with no side chain) is generic and there may be no obvious relation between those labels and labels for real amino acids.

Good-folder sequences are characterized by a large gap $\delta = E_c - E_n$ (compared to the standard deviation σ of the contact energies) between the energy of the designed sequence in the native conformation E_n , and the lowest energy of the conformations structurally dissimilar to the native conformation [1-5]. In other words, good folders are associated with a normalized gap $\xi = \delta/\sigma \gg 1$, a quantity closely related to the *z* score [12]. For example, the 36mer sequence listed in the caption to Fig. 1 and called S_{36} in the literature [13–18], designed by minimizing the energy in the target (native) conformation with respect to the amino acid sequence for fixed composition has, in the units considered here ($RT_{room} = 0.6 \text{ kcal/mol [9]}$), an energy gap $\delta = 2.5$ and thus a sufficiently large value of ξ (=2.5/0.3 \approx 8.33) so as to ensure fast folding. In fact, Monte Carlo (MC) simulations carried out at



FIG. 1. Three conformations are used as natives in the present study. Sequence S_{36} , which is a good folder onto structure (a), is SQKWLERGATRIADGDLPVNGTYFSCKIMENVHPLA.

the temperature T=0.28 of 3000 36mers with energies, in the native conformation, lying inside the gap fold in times $\leq 7 \times 10^7$ MC steps [5] (for caveats see Ref. [19]). In particular S_{36} folds in 6.5×10^6 MC steps.

It has been also shown that most of the thermodynamical [18] and dynamical [5,13] behavior of designed proteins is controlled by only 5-10% of the sites. As a consequence, making mutations in these sites, which are called "hot" in Ref. [18], one destroys, as a rule, the ability the protein has to fold (denaturation). On the other hand, the effects of substitutions in any other site (that can be regarded as "cold") are small, leading to neutral mutations.

III. DESIGNABILITY WITH 20-LETTER MODELS

While 20-letter heteropolymers capture the essential components of real proteins, it is hardly possible to enumerate all sequences which have a given conformation as their nondegenerate ground state. Accordingly, it is not possible to calculate the exact designability of protein conformations. To bypass this problem, we shall determine designability from energetic considerations, using a strategy which relies on the fact that all sequences which have an energy lower than E_c fold on short call, in any case in times which are much shorter than that associated with the random search [20].

Any sequence of a given length N (e.g., N=36) can be obtained making $m \le N$ mutations (i.e., substitution of an amino acid in a given site with a different one) in the minimum energy sequence [e.g., S_{36} in the case of Fig. 1(a)]. Consequently, the designability of a conformation can be found starting from the minimum energy sequence, counting how many mutations lay within the gap $\delta = E_c - E_n$. If ΔE is the change in the energy of the native state produced by a mutation, $p_m(\Delta E)$ the energy distribution probability associated with *m* mutations and n_m^{tot} the total number of sequences that can be produced by introducing *m* mutations in the minimum energy sequence, designability can be defined as

$$n = \sum_{m=1}^{N} n_m, \qquad (2)$$

$$n_m = n_m^{tot} \int_0^\delta p_m(\Delta E) d(\Delta E). \tag{3}$$

So far, we have done nothing more than express the problem in another way, since to know the spectrum of mutation energies of the optimal sequence one has again to enumerate all sequences. In fact, it looks like as if we have made things even worse, in that now one has to find the optimal sequence, which is a nontrivial problem, and also has to ensure that E_c does not change with mutations.

We shall show in the following that the distribution of mutation energies does not depend on the particular structure or on the particular sequence chosen (provided that $E \ll E_c$) nor on the contact energy matrix used to design the protein, but only on its composition and on the number of contacts (or the length, if it is fully compact). This observation leaves room for approximations. In fact, if one is able to find an approximate expression for $p(\Delta E)$, such an expression will hold for all model proteins of the same length. Furthermore, the knowledge of the sequence associated with the global minimum of energy E_n is not necessary (because all sequences have the same spectrum of mutations), only the value of E_n is required. Consequently, even if the optimal sequence cannot be known without a full enumeration of all sequences, it is allowed to use any other sequence with energy $E \approx E_n$, introducing in this way only an error in the integration boundaries (and not on the function to be integrated). It is then possible to calculate designability of a structure from Eqs. (2) and (3) using an approximate distribution $p(\Delta E)$ and an approximate value of δ .

The most conservative way to calculate the number of sequences which fold to a conformation is then to use a distribution $p_m(\Delta E)$ found only by swappings between the residues of the optimal sequence, as in such a way the composition is conserved and E_c does not change. On the other hand, since there are also sequences with different compositions folding to the same conformation, one is also forced, in principle, to calculate the number *n* associated with pointlike mutations. The values found from the swapping of amino acids and from pointlike mutations can be viewed as the lower and the upper limit to designability, respectively.

In Figs. 2(a) and 2(b) we display the unnormalized energy distribution probabilities associated with two compositionconserving and with two pointlike mutations of S_{36} (the integral of these distributions being the total number of sequences). Each of these curves can be well fitted by the sum of two Gaussians, whose means are $\overline{\Delta E_2} = 1.2$ and $\overline{\Delta E_2}$ = 3.0 [Fig. 2(a), composition conserving case] and ΔE_2 = 1.1 and ΔE_2 = 3.6 [Fig. 2(b), pointlike mutations case]. Standard deviations are $\sigma_2 = 0.7$ and $\sigma_2 = 1.0$ [Fig. 2(a)] and $\sigma_2 = 0.7$ and $\sigma_2 = 1.1$ [Fig. 2(b)]. The behavior of these two distributions seems very alike, except for the fact that the area below the composition-conserving curve is much smaller than that below the pointlike mutations curve. This is because much fewer mutations can be made in the first than in the second case and, consequently, the associated Gaussian behavior is less well defined.



FIG. 2. Energy distribution for two composition-conserving (a) and pointlike (b) mutations. The parameters of the Gaussian fit (dotted line) are given in the text.

The overall structure of the curves shown in Fig. 2 can be understood from the fact that the average value of ΔE for "cold" sites is 0.65 and for hot sites is 2.87 [18]. Accordingly, the low-energy peak can be associated with two mutations in cold sites, while the high-energy peak can be associated with a mutation in a cold site and a mutation in a hot site. The contribution from mutations in two hot sites leads to an enhancement of the high-energy tail of the curve. Concerning the Gaussian behavior, we note that the energies associated with the 19 possible mutations on a given cold site are uncorrelated. In other words, one has to pay an energy $\Delta E_2/2 \approx 0.6$ (concerning the factor 1/2 one is reminded of the fact that ΔE_2 gets contributions from two mutations) to remove the wild-type residue, reflecting the fact that it has been optimized. Second, one has to introduce a new residue in the niche left by the wild-type residue. The Gaussian shape of the distribution suggests that the niche is neutral with respect to the new residue and that the new interactions are merely random. To be more precise, the change in energy ΔE upon mutations is the difference between the energy needed to remove the original residue (which is roughly constant and assumes two different values for cold and for hot sites) and the sum of a number of contact energies associated with the new residue, energies which can be considered as random numbers. Being the sum of random numbers, the energy associated with the new residue is forced to respect the central limit theorem, and consequently its distribution approaches a Gaussian function. Of course, an exact Gaussian distribution could be reached only if the number of nearest neighbors of each site were infinite (while in a cubic lattice this is, at most, five). On the other hand, the fact that $p_m(\Delta E)$ can be accurately fitted by Gaussian distributions [cf. e.g., Fig. 2(b)] testifies to the fact that we are not far from the conditions in which the central limit theorem holds.

While the hypothesis that cold mutations give rise to Gaussian-like peaks is quite grounded, due to the uncorrelateness of the energy contributions of cold sites, it is unlikely that the central limit theorem works properly for hot sites, whose energy contributions are correlated [5]. In order to calculate the degree of designability of a protein conformation, we only need to know the contributions from cold sites and, consequently, we do not need to better characterize the peaks associated with hot sites.

We have found that the distribution of mutation energies are rather universal functions. Examples of such a behavior are shown in Fig. 3, where two-pointlike-mutations spectra $p_2(\Delta E)$ associated with low-energy 36mer sequences optimized (making use of the elements of Table 6 of Ref. [8]) on three different conformations [cf. Figs. 1(a) - 1(c)] and with three sequences designed on the same conformation [Fig. 1(a)] are displayed. Similar results have been obtained for chains of different lengths. Furthermore, using different 20 $\times 20$ interaction matrices lead to the same Gaussian behavior of $p_2(\Delta E)$, although the mean values and the standard deviations are different. This is again a consequence of the central limit theorem. This can be seen from Fig. 4, where we display the function $p_2(\Delta E)$ associated with two pointlike mutations on S_{36} (cf. Fig. 1), but making use this time of the interaction matrix elements listed in Table 5 of Ref. [9]. Because, making use of this matrix, the average change in energy upon mutations in cold sites is zero, while that in hot sites is 0.35, it is easy to identify the peaks associated with two cold mutations ($\Delta E_2 = 0$ and $\sigma_2 = 0.34$), with one cold and one hot mutations ($\Delta E_2 = 0.35$ and $\sigma_2 = 0.02$), and with two hot mutations ($\overline{\Delta E_2} = 0.70$ and $\sigma_2 = 0.22$).

Summing up, the function $p_2(\Delta E)$ associated with chains of different length and sequence as well as designed on different native conformations overlap quite nicely, suggesting that the spectrum of both composition conserving and nonconserving mutations is universal. On the other hand, the actual value of $\overline{\Delta E_2}$ and σ_2 characterizing the different peaks of the energy distribution probability depend on the matrix used to describe the contact energies among the amino acids.

The universality of the energy distribution probability is in agreement with the interpretation of the main peaks of the spectrum of mutations of a designed protein in term of cold and hot sites. In fact, the properties of the hot sites are rather homogeneous, their contribution to the mutation spectrum being universal. Assuming, furthermore, that the interactions



FIG. 3. The distribution of energies associated with two composition-conserving mutations made on three sequences designed on three different 36mers conformations. (b) The spectrum obtained making two composition-conserving mutations on three sequences designed on the same conformation (the one displayed in Fig. 2). The values of the energy gap are $\delta = 2.5$ (dotted curve), $\delta = 1.6$ (solid line), and $\delta = 1.3$ (dashed line).

associated with the mutations in cold sites are random, the resulting energy distribution is Gaussian (and so universal), and its standard deviation depends only on the interaction matrix, while its average value depends on the degree of optimization of the wild-type monomer. This, in turn, can be approximated by the degree of optimization of the whole chain (measured by the energy gap δ) divided by its length, a quantity which is essentially constant for long chains [24] (for example, in the case of S_{36} this number is 2.5/36 = 0.07).

IV. HOW MANY GOOD FOLDERS?

The basic idea to calculate the designability of a model protein, as we have discussed above, is to find a simple approximation to the universal distribution of energies associated with mutations onto the optimal sequence, integrate this distribution up to the gap δ , and normalize this result to the total number of mutations that one can make [cf. Eqs. (2) and



FIG. 4. The distribution $p(\Delta E)$ associated with two pointlike mutations for the structure displayed in Fig. 1(a) when the monomers interact with the matrix listed in Table 5 of Ref. [9] (instead of Table 6). The dashed line is the Gaussian fit obtained with the least-squares method.

(3)]. As a consequence of this, designability turns out to depend only on the length of the protein (through the total number of mutations) and on the gap δ .

In order to calculate *n* using Eqs. (2) and (3) we first have to know the total number n_m^{tot} of sequences that can be obtained by making *m* sequence-conserving mutations (swappings) in the optimal sequence. This number can be obtained by counting the number of ways one can select *m* sites, multiplied by the number of permutations of these sites which move all the *m* residues. That is

$$n_m^{tot} = \binom{N}{m} P_0(m), \tag{4}$$

where $P_i(m)$ is the number of ways one can permute *m* sites in such a way that only *i* positions are kept fixed. From the relation

$$m! = P_0(m) + P_1(m) + P_2(m) + \dots + P_{m-2}(m) + 1$$
 (5)

it is possible to extract the expression for P_0 ,

$$P_0(m) = m! - \sum_{k=1}^{m-2} \binom{m}{k} P_0(m-k) - 1.$$
 (6)

For large *m*, one can use the Stirling approximation for the factorials in Eqs. (4)-(6), and keep only the largest exponent term in the sum (saddle point approximation), obtaining $n_m^{tot} \approx \exp(\alpha m)$. The constant α can be determined from the relation $n_N^{tot} = e^{\alpha N} = N!$, which for N = 36 leads to $\alpha = 2.66$.

To proceed further in the calculation of *n*, one needs to find a simple approximation to $p_m(\Delta E)$. For this purpose, we shall express the energy distribution of an arbitrary number of mutations as a convolution of functions $p_2(\Delta E)$ associated with the swapping of two amino acids. The validity of this approximation rests on the ansatz that every couple of mutations affect the energy of the native state independently of the other couple of mutations. This approximation is expected to work also for large values of m, where the probability of mutating neighboring sites is not negligible, because the contact energy associated with the mutated residues are in any case random quantities with average zero (cf. the discussion in the preceding section). Within this scenario, the number of folding sequences displaying 2m mutations and whose energy in the native conformation lies inside the energy gap can be written as

$$n_{2m} \approx n_{2m}^{tot} \int_{0}^{\delta} dE \int_{-\infty}^{+\infty} d\Delta E_{1} d\Delta E_{2} \cdots d\Delta E_{m-1}$$
$$\times p_{2}(\Delta E_{1}) p_{2}(\Delta E_{2}) \cdots p_{2}(\Delta E_{m-1})$$
$$\times p_{2}(\Delta E - \Delta E_{1} - \Delta E_{2} - \cdots - \Delta E_{m-1}).$$
(7)

Making use of the energy distribution probability associated with an amino acid swapping (composition conserving mutations) or with two pointlike mutations (composition nonconserving mutations) one obtains the lower and the upper limit of the designability of a conformation.

In what follows we shall essentially discuss the case of composition conserving mutations. If δ is lower than the peak associated with mutations in hot sites [as in the case of the sequence S_{36} where $\delta = 2.5$, cf. Fig. 2], one should convolute only the peak of $p_2(\Delta E)$ associated with mutations in cold sites [25]. Exploiting the fact that the convolution of *m* Gaussian distributions, of the form $\exp[(\Delta E - \overline{\Delta E_2})^2/2(\sigma_2)^2]$ is a Gaussian distribution with average $\overline{\Delta E_{2m}} = m\overline{\Delta E_2}$ and standard deviation $\sigma_{2m} = m^{1/2}\sigma_2$, it is possible to write

$$n_{2m} \approx n_{2m}^{tot} (2 \pi m^2 \sigma_2^2)^{-1/2} \exp\left(\frac{-(\overline{\Delta E}_2)^2}{2(\sigma_2)^2}m\right)$$
$$\times \int_0^\delta d\Delta E \exp\left(-\frac{\Delta E^2}{2m(\sigma_2)^2} + \frac{\overline{\Delta E}_2 \Delta E}{2(\sigma_2)^2}\right). \quad (8)$$

For $m \ge \delta/(2\sigma_2)^{1/2}$ (in the case of S_{36} this condition means $m \ge 2$) one can neglect the first exponential factor in the integral, in which case the integration can be carried out analytically, leading to

$$n_{m} \approx n_{m}^{tot} (\pi m^{2} \sigma_{2}^{2}/2)^{-1/2} \exp\left(-\frac{\Delta E_{2}^{2}}{4(\sigma_{2})^{2}}m\right)$$
$$\times \frac{2(\sigma_{2})^{2}}{\overline{E}_{2}} \left(\exp\left[\frac{\overline{\Delta E}_{2}}{2(\sigma_{2})^{2}}\delta\right] - 1\right), \tag{9}$$

where the substitution $2m \rightarrow m$ has been made. This equation tells us that designability increases exponentially with the gap δ . In other words, the number of sequences folding to a (compact) conformation is determined only by the gap associated with the minimum energy sequence.

We have shown that the concepts of designability (i.e., the number of sequences folding to a given conformation) and foldability (i.e., thermodynamical stability of the sequences with low energy on the given conformation, expressed by the gap δ) are intimately connected by Eq. (9). If a protein is stable in its native conformation, such native conformation is necessarily highly designable. Vice versa, if a conformation is highly designable, there exist sequences with a large gap folding to it.

To give a numerical evaluation of protein conformations, we make use of Eq. (8), rewritten in the form

$$n = \sum_{m=1}^{N} \frac{k}{m} \exp\left[\left(\alpha - \frac{\overline{\Delta E_2^2}}{4(\sigma_2)^2}\right)m\right], \quad (10)$$

where *k* does not depend on *m* and, for the case of the structure displayed in Fig. 1(a), assumes the value k = 17 (in keeping with the fact that $\delta = 2.5$, $\overline{\Delta E_2} = 1.2$, and $\sigma_2 = 0.7$). In the case in which $\alpha > \overline{\Delta E_2^2/4}(\sigma_2)^2$, which in the case of the 36mer under discussion implies $\alpha > 0.73$, one can keep only the largest term in the above sum. Within this approximation one can write $n \approx e^{1.90 \times 36} = 0.6 \times 10^{30}$, a number to be compared with $n_{36}^{tot} = 3.72 \times 10^{41}$.

One can mention, for the sake of completeness, that the number of sequences within the gap obtained by pointlike mutations (which is the upper limit to designability), is well fitted by the function $4 \exp(5m)$, while the total number of sequences is $19^m {\binom{N}{m}}$.

V. CONCLUSIONS

The degree of designability of a given conformation depends exponentially on the energy gap δ . Since the number of folding sequences is given by the integral of a universal function (the mutation energy distribution) carried up to δ , a quantity which also determines the thermal stability of the designed protein, one can conclude that designability and thermal stability are strongly interconnected. In other words, sequences displaying large gaps are both thermally stable and highly designable. Even sequences displaying, in the native conformation, a small gap fold on short call and share (in the compaction process) the conserved contacts leading to local elementary structures and the (postcritical) folding nucleus [5,26]. Consequently, it is possible to obtain from them, through composition-conserving mutations, other sequences folding to the same native conformation and displaying a large gap. In other words, any sequence able to fold fast, folds to a highly designable conformation.

We have estimated that there are of the order of 10^{30} sequences folding to a compact 36mer conformation, over a total of 10^{41} . This is only the lower limit, but let us assume that it describes well the typical degree of designability of the designed protein. Is this number small or big? The answer to this question has, of course, important implications from the evolutionary point of view. If good folders were distributed homogeneously in the space of sequences (like in the case of RNA [27]) the important parameter would be their density, that is 10^{-11} . This number would be very low, preventing sequences from moving along neutral pathways

(which are collections of sequences folding to the same conformation and differing by single mutations). Such a scenario is very unfavorable for evolution. The situation is, however, quite different for proteins. In fact, it has been shown [28] that good folders group themselves in clusters and superclusters, giving rise to quite an inhomogeneous landscape. Con-

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sequently, the relevant parameter that measures the designability of a conformation is the total number of sequences which conserve, in any way, the energy gap. This number $(>10^{30})$ is very large, in particular in keeping with the fact that over a lifespan of the order of 60 mutations occur in the genome of each person [29].

being equivalent to solving the Fokker-Planck equation for diffusion in a potential, it can be helpful also in studying the kinetical properties of complex systems, provided that the Fokker-Planck approximation is valid (i.e., the moves are local and the potential changes smoothly on the diffusion length scale). Furthermore, Rey and Skolnick have shown [23] that the folding trajectories obtained with Monte Carlo simulations are consistent with those obtained with real molecular dynamics calculations.

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